Stochastic approach to model reduction in computational fluid mechanics. Application to wind time series

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Mascot-Num 2023 Le Croisic

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Announcement



The 2024 annual meeting of MASCOT-NUM will be held in the are a of Nice-Sophia Antipolis, and will be organized by Inria, Université Côte d'Azur LJAD, I3S, EURECOM & Mines ParisTech Sophia.

Save the date!

Local wind variability estimation is relevant in many situations

► For the risk of fatigue evaluation

Pictures from vestas turbines, hub height of 140 m; wind speeds of about 14-18 m/s (ref: from a post on youtube)





► Air quality measurement uncertainty



Need refined short term prediction of wind gust



the Ever-Given into the Suez canal

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• Common features: near wall turbulence; presence of measurement points.

Local wind as a time series - prediction issues

Wind velocity vector measured at a point, at discrete time (with a frequency range from 1 Hz to 50 Hz or more; here t is incremented each 0.1 s.).



 $\langle U_t^{obs} \rangle$ is commonly compute by an average in time over an interval of 10 minutes to 60 minutes, corresponding to a minimum in the wind power spectral density.

Here the mean and intensity are plotted with the time-window $\zeta = 40$ minutes.

Several scales and methods

U(t+k) = U(t)	very short term (seconds to 30 minutes)
Global Forecasting, WRF,	for long term (one day to one week)
ANN, TS-models	for short term (30 minutes to 6 hours)
NWP + ANN,	medium and long term (6 hours to1 week
	$ \begin{array}{l} U(t+k) = U(t) \\ \mbox{Global Forecasting, WRF,} \\ \mbox{ANN, TS-models} \\ \mbox{NWP + ANN,} \end{array} $

[Soman et al., 2010, Chang, 2014, Hanifi et al., 2020].

A double goal

- (1) Propose a Times-series approach, based on SDEs derived from well established physical approaches (that are all including turbulence modelling) to predict the short term distribution of the turbulent velocity.
- (2) Use wind observation as experiments allowing to quantify the uncertainty on supposed well known turbulence modelling parameters.



Fluid particles

taking the perspective of a 'air parcel', and given the flow field $\mathscr{U}(t,x)$, we consider parcel's state variables (x_f, U_f)

1

$$\frac{dx_f}{dt}(t) = U_f(t),$$

$$U_f(t) = \mathscr{U}(t, x_f(t)).$$

But how to get $\mathscr{U}(t, x_f(t))$?



tracers trajectories in turbulence (borrowed from [Bentkamp et al., 2019]).

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2023 is the bicentenary of Navier's work that led to the establishment of the master equations of fluid mechanics, known as the **Navier-Stokes equations** that governing $\mathscr{U}(t, x)$

$$\partial_t \mathscr{U}^{(i)} + \mathscr{U}^{(j)} \partial_{x_j} \mathscr{U}^{(i)} = \nu \triangle \mathscr{U}^{(i)} - \frac{1}{\varrho} \partial_{x_i} \mathscr{P}$$
$$\partial_{x_i} \mathscr{U}^{(i)} = 0$$

Direct Numerical Simulation are from very expensive to totally prohibitive, as it requires a mesh below the Kolmogorov length scale η_{K} in [50µm, 1mm] for most of industrial or environmental flows.

Averaged Navier Stokes equations

 $\mathscr{U}(t,x) \rightsquigarrow \langle \mathscr{U} \rangle(t,x) + {\rm a}$ model for the 2nd moments lost with the subscales

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for CFD.

 $\blacktriangleright \mathscr{U}(t,x) = \langle \mathscr{U} \rangle(t,x) + \text{noise.}$ Lagrangian modelling requires a model for the noise

$$\frac{dx_f}{dt}(t) = \langle \mathscr{U} \rangle_{\text{Ens, LES, RANS}}(t, x_f(t)) + u(t)$$

with u(t) a random fluctuation of the Lagrangian mean velocity (Lagrangian Particle Dispersion Model (LPDM)).

Turbulent second order closure (See e.g. [Durbin and Speziale, 1994, Pope, 1994]).

Macroscopic random fluctuation, assuming decorrelation of time increments that lead to Gaussian fluctuation and 3D-Brownian motion B:

$$du^{(i)}(t) = -\frac{u^{(i)}(t)}{\tau_f} dt + (C_0 \varepsilon)^{1/2} dB_t^{(i)}, \qquad (\text{the simplest Langevin model})$$

 C_0 Kolmogorov constant and ε the dissipation rate of the mean kinetic energy required.

Stand alone Lagrangian stochastic model

Modelling consistency: the conditional mean field of air parcel velocity is the conditional expectation of its velocity

$$\langle U_f \rangle(t,x) = \mathbb{E} \Big[U_f(t) \underbrace{ | x_f(t) = x }_{\text{conditionning}} \Big]$$

on $(\Omega, \mathcal{F}, \mathbb{P}, B)$, with (x_f, U_f) solution of a General Langevin Model:

 $\begin{aligned} dx_f(t) &= U_f(t) \, dt, \\ dU_f^{(i)}(t) &= -\partial_{x_i} \langle \mathscr{P} \rangle(t, x_f(t)) dt + \left(G_{ij} \left(U_f^{(j)} - \langle U_f^{(j)} \rangle \right) \right)(t, x_f(t)) dt + \sigma_{ij}(t, x_f(t)) dB_t^{(i)} \end{aligned}$

B is a 3D-Brownian motion. (see e.g. [Durbin and Speziale, 1994, Pope, 2000, Minier and Peirano, 2001])

$$G_{ij} = -\frac{C_R}{2} \frac{\varepsilon}{\mathbf{k}} \delta_{ij} + C_2 \partial_j \langle U_f^{(i)} \rangle, \qquad \sigma_{ij} = \frac{2}{3} \left(C_R \varepsilon + C_2 \mathcal{P} - \varepsilon \right) \delta_{ij},$$

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▶ The probabilistic model is a McKean Vlasov SDE, with

$$egin{aligned} \langle \cdot
angle(t,x) &= \mathbb{E}[\cdot | x_f(t) = x] \ u(t) &= U_f(t) - \langle U_f
angle(t,x_f(t)) \end{aligned}$$

 $\mathcal{P} = \frac{1}{2} \mathcal{P}_{ii}, \text{ the turbulent production term } \mathcal{P}_{ij} := -(\langle u^{(i)} u^{(k)} \rangle \partial_k \langle U_f^{(i)} \rangle - \langle u^{(j)} u^{(k)} \rangle \partial_k \langle U_f^{(j)} \rangle, \\ \varepsilon \text{ is closed with coherent parametrisation involving } \mathbf{k} = \frac{1}{2} \langle u^{(i)} u^{(i)} \rangle.$

Numerical Stand alone Lagrangian stochastic model (in-house SDM code)

Almeida - Particle Trayectory - 2D (X,Z) - t=0.73s



[Mokrani et al., 2019]

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Very hight resolution simulation (downscaled from WRF)



Left: Numerical domain of the water body in WRF+SDM-WindPoS. Synchronised snapshot of the wind magnitude during the day on 24 April 2021 at the first height (10 m) of SDM-WindPoS. Middle: the subdomain is resolved to 150 m. Right: the resolution is 50 m.

Collaboration with SportRizer & Risk Weather Tech

From 3D+time Averaged Navier-Stokes equations to reduced 0D+time SDE

From the Generalized Langevin Model (GLM),

$$dx_{f}^{(i)}(t) = U_{f}^{(i)}(t)dt, \ 1 \le i \le 3, \\ dU_{f}^{(i)}(t) = -\frac{1}{\rho}\partial_{i}\langle\mathscr{P}\rangle(t, x_{f}(t)) \ dt + G_{ij}(t, x_{f}(t)) \underbrace{(U_{f}^{(j)}(t) - \langle U_{f}^{(j)}\rangle(t, x_{f}(t)))}_{u(t) \ \text{turb. velocity}} dt + (C_{0}\varepsilon)^{1/2} (t, x_{f}(t)) \ dB_{t}^{(i)}$$
(1)

 $\langle \mathscr{P} \rangle$ is the mean pressure

Fix $x_f(t) = x_{obs}$,

$$u_t = U_f(t) - \langle U_f \rangle(t, x_{\text{obs}}) \qquad \langle U_f^{(i)} \rangle(t, x_{\text{obs}}) = \mathbb{E}[U_f^{(i)}(t)|x_f(t) = x_{\text{obs}}],$$

Then the SDE for the instantaneous turbulent velocity $(u_t, t \ge 0)$ seen at x_{obs} is

$$du_t^{(i)} = G_{ij}(t, x_{\text{obs}})u_t^{(j)}dt + \sqrt{C_0(t, x_{\text{obs}})\varepsilon(t, x_{\text{obs}})} \, dB_t^{(i)},$$

and its squared norm

$$d||u_t||^2 = -2u_t^{(j)}G_{ij}(t, x_{\text{obs}})u_t^{(j)}dt + 3(C_0\varepsilon)(t, x_{\text{obs}})dt + 2\sqrt{(C_0\varepsilon)(t, x_{\text{obs}})}||u_t|| dW_t,$$

with the process $W_t = \sum_i \int_0^t \frac{u_s^{(i)}}{\|u_s\|} dB_s^{(i)}$ identifies as a one-dimensional Brownian motion.

• We choose the isotropic (diagonal) structure for G_{ij}

$$G_{ij}(t,x) = -\frac{C_R}{2} \frac{\varepsilon}{\mathbf{k}}(t,x) \,\delta_{ij}, \qquad \text{where} \quad \mathbf{k}(t,x) = \frac{1}{2} \mathbb{E}[\|u_t\|^2 |x_f(t) = x] \tag{2}$$

corresponding to the Simplified Langevin model [Pope, 1985]. For consistency reason, C_0 is now $C_0=\frac{2}{3}(C_R-1).$

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• We choose a model for the dissipation rate of the kinetic energy $\varepsilon(t,x)$, classically used in the atmospheric boundary layer (ABL) :

$$\varepsilon(t,x) = \frac{C_{\varepsilon}}{\ell_{\rm m}} \mathbf{k}^{3/2}(t,x), \tag{3}$$

where C_{ε} is a constant, $\ell_{\rm m}$ is a characteristic length scale called *m*ixing length. Near the ground, $\ell_{\rm m} = \kappa z$, where *z* denotes the distance to the wall from *x*, and κ is the Von Kármán constant (Cuxart et al., 2000, Drobinski et al., 2006). • We choose the isotropic (diagonal) structure for G_{ij}

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 C_0, C_R might vary according to the model and context : $C_R \in [1.5, 5]$, implying $C_0 \in [\frac{1}{3}, \frac{8}{3}]$ We set $C_\alpha = \frac{C_{\varepsilon}}{\kappa z(x_{obs})}$, with $z(x_{obs}) = 30$ m. Von Kármán constant $\kappa \in [0.287, 0.615]$ [Edeling et al., 2014].

 $C_{\mu} \in [0.054, 0.135]$ [Edeling et al., 2014], for $C_{\varepsilon} = C_{\mu}^{3/4}.$

 \triangle We expect the values of C_{α} to be within the interval [0.0061, 0.0259].

The reduced **0D+time SDEs**

At x_{obs} , with notation $q_t = ||u_t||^2$, such that $k(t, x_{obs}) = \frac{1}{2}\mathbb{E}[q_t]$, incorporating both $G_{ij} = -\frac{C_R}{2}\frac{\varepsilon}{k}\delta_{ij}$ and the parametrisation of the dissipation $\varepsilon = C_{\alpha}k^{\frac{3}{2}}$, we end up with our physical-based model obtaining the following **CIR-type stochastic mean-field TKE** model:

$$dq_{t} = \gamma dt - C_{R} \frac{C_{\alpha}}{\sqrt{2}} q_{t} (\mathbb{E}[q_{t}])^{\frac{1}{2}} dt + 3C_{0} \frac{C_{\alpha}}{2\sqrt{2}} (\mathbb{E}[q_{t}])^{\frac{3}{2}} dt + \sqrt{\sqrt{2}C_{0}C_{\alpha}} (\mathbb{E}[q_{t}])^{\frac{3}{4}} \sqrt{q_{t}} dW_{t}, \quad (4)$$

 $C_0 = \frac{2}{3}(C_R - 1); q_0$ given.

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[Bossy et al., 2022]

- (1) For positive parameters C_{α} , C_0 , and γ , pathwise wellposedness holds for $(q_t; t \ge 0)$ solution of the McKean-Vlasov SDE (4).
- $(2) \ \underline{C} \ \mathbb{1}_{\gamma>0} \leq \sup_{t\geq 0} \mathbb{E}[q_t^p] \leq \overline{C}, \text{ for all } p\geq 1, \text{ for some constants } \overline{C}, \underline{C}>0.$

(3) For p = 1,

$$q_0 \wedge \left(\frac{\sqrt{2\gamma}}{C_{\alpha}}\right)^{2/3} \leq \sup_{t \geq 0} \mathbb{E}[q_t] \leq q_0 \vee \left(\frac{\sqrt{2\gamma}}{C_{\alpha}}\right)^{2/3},$$

with long-time behaviour

$$\lim_{t \to +\infty} \mathbb{E}[q_t] = \left(\frac{\sqrt{2\gamma}}{C_{\alpha}}\right)^{2/3}$$

using wellposedness condition for CIR processes with time dependant coefficients [Maghsoodi, 1996].

A two steps strategy

- 1. Construct first a simple maximum likelihood estimator to recover the model constants from the data.
- 2. Next introduce a Bayesian estimation to quantify the uncertainty on theses parameters.

Difficulty : McKean Vlasov form of the model.

Two linearization tentatives :

• $\mathbb{E}[q_t] \rightsquigarrow q_t$

$$dq_t = \gamma dt - Aq_t^{3/2} dt + Bq_t^{5/4} dW_t$$

• $\mathbb{E}[q_t] \rightsquigarrow \mathbb{E}[q_\infty]$, explicitly known, leading to

$$dq_t = \Theta \left(\mu - q_t\right) dt + \sigma \sqrt{q_t} dW_t,$$

Linearisation $\mathbb{E}[q_t] \longrightarrow q_t$: Wellposedness and control of moments

$$dq_t = \gamma dt - \frac{C_{\alpha}}{\sqrt{2}} q_t^{3/2} dt + \sqrt{\sqrt{2}C_0 C_{\alpha}} q_t^{5/4} dW_t, \quad q_0 = |u_0'|^2$$

SDEs with superlinear growth coefficients :

$$dY_t = \gamma dt - B Y_t^{2\alpha - 1} dt + \sigma Y_t^{\alpha} dW_t, \ Y_0 = y_0 > 0, \quad \text{with } \alpha > 1.$$

[Bossy et al., 2022]

Assume $\gamma, B \ge 0$. Then there exists a unique (strictly) positive strong solution Y to the SDE (5).

 $\begin{array}{l} \text{Moments: For } p \text{ such that } 0 \leq 2p \leq 1 + \frac{2B}{\sigma^2}, \qquad \sup_{t \in [0,T]} \mathbb{E}\big[Y_t^{2p}\big] < +\infty. \\ \text{Exponential moments: for all } \mu \leq C(\alpha, \gamma), \sup_{t \in [0,T]} \mathbb{E}\big[\exp\{\mu \int_0^t Y_s^{2\alpha-2} ds\}\big] < +\infty. \end{array}$

A Here $\frac{2B}{\sigma^2} = \frac{1}{C_0}$, so only $C_0 \in [\frac{1}{3}, 1]$ ensures that $\mathbb{E}[q_t^2]$ is finite.

- The (linear) model behave in time like the McKean Vlasov model (when $\gamma = 0$) $\lim_{t \to +\infty} \mathbb{E}[q_t] = 0, \quad \lim_{t \to +\infty} t^2 \mathbb{E}[q_t] = \frac{4}{C_{\alpha}^2(C_0^2 + 3C_0 + 2)}.$
- The Euler scheme $q_{t_{n+1}} = q_{t_n} + B|q_{t_n}|^{3/2} + \sigma|q_{t_n}|^{5/4}(W_{t_{n+1}} W_{t_n})$ is almost surely converging (Gyongy, 1988), allowing to propose a consistent Maximun Likelyhood estimator for $(C_{\alpha}, C_0, \gamma)$.

▲ Falling in case of strong L¹ divergence of the Euler scheme [Hutzenthaler et al., 2010] for SDEs with super-linear growth condition.

(5)

Linearisation
$$\mathbb{E}[q_t] \longrightarrow \mathbb{E}[q_\infty] = \left(rac{\sqrt{2}\gamma}{C_lpha}
ight)^{2/3}$$

CIR model for the instantaneous TKE:

$$dq_t = \Theta(C_\alpha, \gamma) \left(\mu(C_\alpha, \gamma) - q_t \right) dt + \sigma(\gamma) \sqrt{q_t} dW_t, \quad q_0 \text{ given, and } \gamma > 0, \tag{6}$$

where

$$\Theta(C_{\alpha},\gamma) = C_R(\frac{C_{\alpha}^2\gamma}{2})^{1/3}, \quad \mu(C_{\alpha},\gamma) = (\sqrt{2}\frac{\gamma}{C_{\alpha}})^{2/3}, \quad \sigma(\gamma) = \sqrt{2C_0\gamma}$$

with physical parameters : $(\gamma, C_0, (C_R), C_\alpha)$

 $\mathbb{P}(\inf\{t \ge 0, q_t = 0\} = +\infty) = 1 \Leftrightarrow 2\Theta(C_\alpha, \gamma)\mu(C_\alpha, \gamma) \ge \sigma^2(\gamma) \Leftrightarrow C_R \ge C_0 \text{ structurally always satisfied.}$

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Symmetrized Euler scheme : for SDE (6), $t_n = n\Delta t$

$$\widehat{q}_{t_{n+1}} = \left| \widehat{q}_{t_n} + \Theta(C_\alpha, \gamma) \left(\mu(C_\alpha, \gamma) - \widehat{q}_{t_n} \right) \Delta t + \sigma(\gamma) \sqrt{\widehat{q}_{t_n}} \left(W_{t_{n+1}} - W_{t_n} \right) \right| \tag{7}$$

Equivalently

$$\widehat{q}_{t_{n+1}} \sim |\mathcal{N}(\widehat{q}_{t_n} + \Theta(C_\alpha, \gamma)(\mu(C_\alpha, \gamma) - \widehat{q}_{t_n})\Delta t, \ \sigma^2(\gamma)\widehat{q}_{t_n}\Delta t)|.$$

Under the assumption $C_0 < 2$, the scheme (7) converges in law with a rate one [Bossy and Diop. 2010]. We fix $C_0 = 1.9$ (in accordance with the literature).

$(\widehat{C}_{lpha},\widehat{\gamma})$ estimators

Considering the compact set $D \subset \mathbb{R}^+ \times \mathbb{R}^+$ supporting the admissible values of $\theta = (C_{\alpha}, \gamma)$, choosing Δt according to some data frequency, we compute the pseudo-maximum likelihood estimator

$$\widehat{\theta} = \underset{\theta \in D}{\arg \max} \ \log p^{\theta}_{\Delta t}(q_0^{\mathsf{obs}}, \dots, q_{t_N}^{\mathsf{obs}}),$$

allowing for an explicit solution of the optimal pair $(\widehat{C}_{\alpha}, \widehat{\gamma})$

Define empirical moments

$$\widehat{M}_{m_1,m_2} = \frac{1}{N} \sum_{n=0}^{N-1} (q_{t_n+1}^{\text{obs}} - |q_{t_n}^{\text{obs}}|)^{m_1} |q_{t_n}^{\text{obs}}|^{m_2}.$$
(8)

Quadratic variation estimator for γ :

$$\widehat{\gamma} = \frac{\widehat{M}_{2,0}}{2C_0 \Delta t \widehat{M}_{0,1}},\tag{9}$$

Pseudo-maximum likelihood estimator of C_{α} :

$$\widehat{C}_{\alpha} = \frac{\sqrt{2}}{\sqrt{\widehat{\gamma}}} \max\left\{ c_* \frac{\sqrt{\widehat{\gamma}}}{\sqrt{2}}, \left(\frac{\max\{\widehat{\gamma} \Delta t C_R - \widehat{M}_{1,0}, 0\}}{\widehat{M}_{0,1} \Delta t C_R} \right)^{3/2} \right\}$$
(10)

for some lower bound c_* to choose in [0, 0.0061].

We end up with a family of estimators

 $\Sigma := \{ (\widehat{C}_{\alpha}(d), \ \widehat{\gamma}(d)), \text{ for } d \text{ in the selection of day-periods } \$ \text{ in the year 2017} \},$



Point-estimations for each Wednesday of 2017.

The dark grey area (right) is the reference interval compiling turbulence closure literature.

Priori distributions, defined as truncated Gaussian distributions:

$$\gamma \sim \mathcal{N}^+ \left(\overline{\Gamma}(\mathcal{S}), \mathbb{V}_{\Gamma}(\mathcal{S})\right), \qquad C_{\alpha} \sim \mathcal{N}^+ \left(\overline{C}(\mathcal{S}), \mathbb{V}_C(\mathcal{S})\right),$$

with $\overline{\Gamma}(\$), \overline{C}(\$), \mathbb{V}_{\Gamma}(\$)$ and $\mathbb{V}_{C}(\$)$ are the empirical means and variances over the days.

Step one: Posterior calibration

Goal : use Bayesian inference to extract precise information on the parameters distributions $\pi(\theta|q^{\text{obs}})$ from the data q^{obs} , with $\theta = (\gamma, C_{\alpha})$,

$$\pi(\theta|\boldsymbol{q}^{\text{obs}}) = \frac{p(\boldsymbol{q}^{\text{obs}}|\theta) p_{\theta}(\theta)}{p(\boldsymbol{q}^{\text{obs}})},$$
(11)

where

- $p(\cdot | \theta)$ the probability density of the model given the parameters (likelihood function, given)
- p_{θ} the prior distribution of θ (prior distribution, given),
- + $p(q^{\text{obs}})$ is the distribution of the observed data with

$$q^{\mathrm{obs}}(\theta) = \ \widehat{q}(\theta) + \mathscr{C}$$

- $\hat{q}(\theta)$ is the i.i.d random vector variable, with the equilibrium law (for a given θ) of the discrete-time model (7);
- the random Error vector $\mathscr{C} \sim$ centred logistic distribution and scale parameter to be estimated from the data. Choice made from a step0 analysis fitting the histogram of observation error distribution with a set of explicit mean-variance probability densities.



Method : Metropolis-Hasting algorithm and its Hamiltonian Monte Carlo (HMC) variant.

We have used the Python package PyMC3 [Salvatier et al., 2016] with the No U-Turn sampler[Gelman and Hoffman, 2014] method.

Posterior calibration on C_{α}



Box plots within month of posterior distribution of \widehat{C}_{α} , constructed from the Markov chains samples in a given month.



Two examples of Exploration of the state space with the Markov chain and posterior distribution

Bayesian calibration of γ during February, 2017

Box plot of the $\gamma(d, i)$ for $i = 1, \dots, 48$ for each 20 minutes-length sub-signal

The $t \mapsto \overline{\gamma}_t$ obtained from Step one for the same four days; the horizontal lines are the level of the means over the period (the black line is the Step zero estimator).



Validation of the calibration procedure against the observation



Instantaneous turbulent kinetic energy observed during February, 2017, between 5 a.m and 8 p.m (color plots) using the frequency of 1/30 s^{-1}

95% confidence interval (plotted in black) of (q_{t_n}, n) using the posterior mean values $\widehat{C}_{\alpha}(d)$ and the time dependent mean $\overline{\gamma}(t) = \sum_{i=0}^{S} \mathbb{E}[\gamma(d, i)] \mathbb{1}_{[T_i, T_{i+1}]}(t)$, with $\Delta t = 30s$.

CIR model for the instantaneous TKE:

$$dq_t = \Theta(C_\alpha, \gamma(t))(\mu(C_\alpha, \gamma(t)) - q_t) dt + \sigma(\gamma(t))\sqrt{q_t} dW_t$$

suggests that $\gamma(t) - \left(\frac{C_{\alpha}^2 \gamma(t)}{2}\right)^{1/3} \mathbb{E}[q_t] = 0$, from which we deduce the formal relation:

$$\gamma(t) = \frac{C_{\alpha}}{\sqrt{2}} \left(\sqrt{3} \| \langle U_{(d)}^{\text{obs}} \rangle \| I_t \right)^3.$$
(12)

Putting the empirical $I_t = \sqrt{\langle ||U_t^{\text{obs}} - \langle U_t^{\text{obs}} \rangle ||^2 \rangle} / \sqrt{3} ||\langle U_{(d)}^{\text{obs}} \rangle|| \text{ in } (q_{t_n}, n).$



95% Confidence interval of (q_{t_n}, n) obtained using the posterior mean values $\hat{C}_{\alpha}(d)$ and the time dependent mean $\bar{\gamma}(t) = \sum_{i=0}^{S} \mathbb{E}[\gamma(d, i)] \mathbb{1}_{[T_i, T_{i+1}]}(t), \Delta t = 30s.$



Prediction of the 95% Confidence interval of (q_{t_n}, n) obtained by sampling the within-year posterior distribution of C_{α} and the time dependent mean $\bar{\gamma}(t)$ replaced by the turbulent intensity statistic through (12) $\Delta t = 30 s$.

In green: the instantaneous turbulent kinetic energy observed during February 15th, 2017, between 7 a.m and 8 p.m using the frequency of 1/30 s.

As a conclusion (work in progress) : Gust modelling and intermitency

The wind gust speed $U_{\rm max}$ is defined as a short-duration maximum of the horizontal wind speed [Suomi and Vihma, 2018]

$$U_{\max}(\overline{t}) = \max\left\{U_f(s); s \in [\overline{t} - h^g, \overline{t}]\right\}.$$

(the choice of the gust duration h^g may vary with the activity sector).

Some predictive frameworks are ready to use, for example

 $U_{\max}(\bar{t}) = \langle U \rangle(\bar{t}) + \text{turbulent kinetic energy} \times g_x.$

where, for a R level is fixed, the peak factor g_x , such that

 $\mathbb{P}(U_{\max}(\bar{t}) < g_x) = 1 - R$

is given by some model formula (under assumption of stationary, and Gaussian behaviour of the acceleration) [Schreur and Geertsema, 2008].



within the CIR model

For a significance level R, we want to compute the probability

$$\mathbb{P}(q_t + \Delta t > g_x | q_t) = R$$

- Improvement of the model and noises
 - · Introduce acceleration :

$$\begin{split} dx_f(t) &= U_f(t) \, dt, \\ dU_f^{(i)}(t) &= -\partial_{x_i} \langle \mathscr{P} \rangle(t, x_f(t)) dt + \left(G_{ij} \left(U_f^{(j)} - \langle U_f^{(j)} \rangle \right) \right)(t, x_f(t)) dt + a(t) dt \\ da(t) &= -\beta a(t) dt + \sigma_{ij}(t) dB_t^{(i)} \end{split}$$

[Innocenti et al., 2020]

- Introduce intermittency : K62, the dissipation ε is a log-normal process. $\varepsilon_t = \mathbb{E}[\varepsilon_t] \exp(\sqrt{\lambda_I} X_t + \frac{\lambda_I}{2} \mathbb{E}[X_t^2])$ with X_t
 - Ornstein-Uhlenbeck (K62)
 - Stochastic Voltera equation reaching log correlation [Letournel et al., 2021]

Some concluding remarks

Modelling of the distribution of the instantaneous wind speed for short term prediction purpose.

- Starting from the physical turbulent modelling, a 3d+time velocity field dynamics, we end up with several propositions of 0d+time stochastic models.
- The simplest proposition (CIR TKE)
 - allows to recover from measurements parameters values that fit the interval values given by the expert's knowledge, as a validation of the physical meaning of the model.



allows to give a good prediction of the 95% CI

- · Still some need for stable explicit schemes for MLE prior calibration purpose
- Raising the still challenging guestion about MLE for McKean Vlasov
- Improving the modelling thorough uncertainty parameters inference.

Thank you for your attention !

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